The Time-Dependent Hartree-Fock-Bogoliubov Equations for Bosons

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## Bose-Einstein Condensate

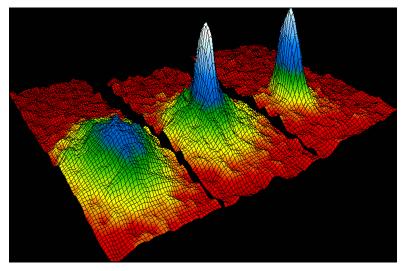


Figure : Velocity-distribution for a gas of rubidium atoms, showing the transition to Bose-Einstein condensation (1995).

## Motivation

- Understand the dynamics of a Bose-Einstein condensate.
- In particular the quantum fluctuations with respect to the evolution through the time-dependant Gross-Pitaevskii equation.

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#### Model for Many Bosons

• Hamiltonian on  $L^2_s(\mathbb{R}^{dN})$  of the form

$$H_N = \sum_{n=1}^N -\Delta_{x_n} + \sum_{1 \le n \le m \le N} v(x_m - x_n)$$

An approximation is

$$e^{-itH_N}\phi_0^{\otimes N}\simeq \phi_t^{\otimes N}$$

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with  $\phi_t$  solution to the Hartree equation:

$$i\partial_t\phi_t = -\Delta\phi_t + (\mathbf{v}*|\phi_t|^2)\phi_t$$

## Model for Many Bosons

• Hamiltonian on  $\mathcal{F} = \bigoplus_{N \ge 0} L_s^2(\mathbb{R}^{dN})$  of the form

$$\mathbb{H} = \bigoplus_{N \ge 0} H_N = \bigoplus_{N \ge 0} \left( \sum_{n=1}^N -\Delta_{x_n} + \sum_{1 \le n < m \le N} v(x_m - x_n) \right)$$

An approximation is

$$e^{-it\mathbb{H}}E(\phi_0) = \bigoplus_{N\geq 0} \frac{e^{-\frac{\|\phi_0\|^2}{2}}}{\sqrt{N!}} e^{-itH_N} \phi_0^{\otimes N} \simeq \bigoplus_{N\geq 0} \phi_t^{\otimes N} \frac{e^{-\frac{\|\phi_0\|^2}{2}}}{\sqrt{N!}} = E(\phi_t)$$

with  $\phi_t$  solution to the Hartree equation:

$$i\partial_t \phi_t = -\Delta \phi_t + (\mathbf{v} * |\phi_t|^2) \phi_t$$

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## **Quasifree States**

The definition of Quasifree states is technical: omitted. Key facts about quasifree states:

States on the Fock space,

i.e. positive linear forms  $\omega^{qf}$  on the bounded linear operators on  $\mathcal{F}(L^2(\mathbb{R}^d)) = \bigoplus_{N=0}^{\infty} L^2_s(\mathbb{R}^{dN})$ , s. t.  $\omega^{qf}(Id_{\mathcal{F}}) = 1$ .

 If L<sup>2</sup>(ℝ<sup>d</sup>) is replaced by C, then F(C) ≅ Barg(C) ⊂ L<sup>2</sup>(C), and quasifree states are "gaussians" in L<sup>2</sup>(C).

Characterized by their (truncated) moments of order 1 and 2.

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 $\begin{array}{ll} (\text{Order 1}) & \phi \in L^2(\mathbb{R}^d), \\ (\text{Order 2}) & \gamma \in \mathfrak{S}^1(L^2(\mathbb{R}^d)), \ \gamma^* = \gamma \ge 0, \\ (\text{Order 2}) & \sigma \in L^2(\mathbb{R}^{2d}) \simeq \mathfrak{S}^2(L^2(\mathbb{R}^d)), \ \sigma^T = \sigma. \end{array}$ 

#### Equivalence of the three dynamics

Theorem (Bach, B., Chen, Fröhlich, Sigal, 2016) Assume  $v^2 \leq 1 - \Delta$ , v(x) = v(-x), and, with  $M = (1 - \Delta)^{1/2}$ ,  $\phi_t \in H^1(\mathbb{R}^d)$ ,  $\gamma_t \in M^{-1} \mathfrak{S}^1(L^2(\mathbb{R}^d)) M^{-1}$ ,  $\sigma_t \in H^1(\mathbb{R}^{2d})$ ,  $\gamma^* = \gamma \geq 0$ ,  $\sigma^T = \sigma$ . The we have equivalence of the three dynamics:

1. For all observable  $\mathbb{A}$  at most quadratic in the fields,

$$i\partial_t \omega_t^{qf}(\mathbb{A}) = \omega_t^{qf}([\mathbb{A},\mathbb{H}])$$

(φ<sub>t</sub>, γ<sub>t</sub>, σ<sub>t</sub>) satisfy the HFB equations (given in the next slide)
 For all observable A

$$i\partial_t \omega_t^{qf}(\mathbb{A}) = \omega_t^{qf}\left([\mathbb{A}, \mathbb{H}_{HFB}^{(2)}(\omega_t^{qf})]\right),$$

with  $\mathbb{H}_{HFB}^{(2)}(\omega_t^{qf}) = \cdots$ .

#### Hartree-Fock-Bogoliubov Equations for Bosons

Hartree-Fock-Bogoliubov Equations for Bosons:  

$$i\partial_t \phi_t = \mathbf{h}(\gamma_t)\phi_t + k(\sigma_t)\bar{\phi}_t + (\mathbf{v} * |\phi_t|^2)\phi_t ,$$

$$i\partial_t \gamma_t = [\mathbf{h}(\gamma_t^{\phi_t}),\gamma_t]_- + [k(\sigma_t^{\phi_t}),\sigma_t]_- ,$$

$$i\partial_t \sigma_t = [\mathbf{h}(\gamma_t^{\phi_t}),\sigma_t]_+ + [k(\sigma_t^{\phi_t}),\gamma_t]_+ + k(\sigma_t^{\phi_t}) ,$$

where  $(v \sharp A)(x; y) := v(x - y)A(x; y)$ ,  $A \in \mathcal{B}(L^2(\mathbb{R}^d))$  and

$$\begin{split} h(\gamma) &:= -\Delta + \mathbf{v} * \gamma(\mathbf{x}; \mathbf{x}) + \mathbf{v} \sharp \gamma , \qquad k(\sigma) := \mathbf{v} \sharp \sigma , \\ \gamma_t^{\phi_t} &:= \gamma_t + |\phi_t\rangle \langle \phi_t| , \qquad \qquad \sigma_t^{\phi_t} := \sigma_t + \phi_t \otimes \phi_t , \\ [\mathbf{A}, \mathbf{B}]_- &:= \mathbf{A} \mathbf{B}^* - \mathbf{B} \mathbf{A}^* , \qquad \qquad [\mathbf{A}, \mathbf{B}]_+ := \mathbf{A} \mathbf{B}^T + \mathbf{B} \mathbf{A}^T . \end{split}$$

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## Existence and Uniqueness

Theorem (Bach, B., Chen, Fröhlich, Sigal, 2016) Assume  $v^2 \leq 1 - \Delta$ , v(x) = v(-x), and  $\phi_0 \in H^1(\mathbb{R}^d)$ ,  $\gamma_0 \in M^{-1} \mathfrak{S}^1(L^2(\mathbb{R}^d)) M^{-1}$ ,  $\sigma_0 \in H^1(\mathbb{R}^{2d})$ , Then there is existence and uniqueness of maximal mild solutions  $(\phi_t, \gamma_t, \sigma_t)$  to the HFB equations. If moreover

$$egin{pmatrix} \gamma_0 & \sigma_0 \ \sigma_0^* & 1+ar{\gamma}_0 \end{pmatrix} \geq 0\,,$$

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then the solution is global.

Theorem (Bach, B., Chen, Fröhlich, Sigal, 2016) Assume  $v^2 \leq 1 - \Delta$ , v(x) = v(-x), and  $\phi_t \in H^1(\mathbb{R}^d)$ ,  $\gamma_t \in M^{-1} \mathfrak{S}^1(L^2(\mathbb{R}^d)) M^{-1}$ ,  $\sigma_t \in H^1(\mathbb{R}^{2d})$ . Then

- the energy  $\omega_t^{qt}(\mathbb{H})$  is preserved by the HFB equations,
- the number of particles is preserved,
- For any obervable A at most quadratic in the field and such that [A, ℍ] = 0, ω<sub>t</sub><sup>qf</sup>(A) is preserved.

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### **Previous Works**

- 1996 Griffin. Conserving and gapless approximations for an inhomogeneous Bose gas at finite temperatures.
- 1998 Dodd, Edwards, Clark, and Burnett. *Collective excitations of Bose-Einstein-condensed gases at finite temperatures.*
- 1998 Parkins and Walls. *The physics of trapped dilute-gas Bose–Einstein condensates.*
- 2010/2011 Grillakis, Machedon, and Margetis. Second-order corrections to mean field evolution of weakly interacting bosons. I & II.
- 2015/arxiv Grillakis and Machedon. Pair excitations and the mean field approximation of interacting bosons, I & II.
  - 2015 Lewin, Nam, and Schlein. *Fluctuations around Hartree states in the mean-field regime.*
  - arxiv Nam and Napiórkowski. Bogoliubov correction to the mean-field dynamics of interacting bosons.
  - arxiv Lewin. Mean-field limit of bose systems: rigorous results.
  - arxiv Boccato, Cenatiempo and Schlein. Quantum many-body fluctuations around nonlinear Schrödinger dynamics.

# Thank you for your attention.

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